# Vortex Reconnections across the BCS-BEC Crossover

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### **1. Density Functional Theory for Fermi Superfluids**

#### **Bogoliubov-de Gennes equations**

- 1. The energy functional of the mean-field superfluid Fermi gas in the BCS regime is  $E = \int d\mathbf{x} \, \mathcal{E}(\mathbf{x}), \quad \mathcal{E}(\mathbf{x}) = \hbar^2 (\tau_{\uparrow} + \tau_{\downarrow})/2m + |\Delta|^2/g$ 
  - $\tau_{\uparrow} = \sum |\nabla u_{\eta}|^2 f(E_{\eta}), \quad \tau_{\downarrow} = \sum |\nabla v_{\eta}|^2 f(-E_{\eta}), \quad \Delta(\mathbf{x}) = -g_{\text{eff}} \sum u_{\eta}(\mathbf{x}) v_{\eta}^*(\mathbf{x})$
- 2. The energies (  $E_{\eta}$  ) and wave functions of quasiparticles  $(u_{\eta}, v_{\eta})^T$  are obtained as self-consistent solutions to the Bogoliubov-de Gennes equations  $\begin{pmatrix} \mathcal{H}_{BdG} - \mu_{\uparrow} & \Delta(\mathbf{x}) \\ \Delta^{\dagger}(\mathbf{x}) & -\mathcal{H}_{BdG} + \mu_{\downarrow} \end{pmatrix} \begin{pmatrix} u_{\eta}(\mathbf{x}) \\ v_{\eta}(\mathbf{x}) \end{pmatrix} = E_{\eta} \begin{pmatrix} u_{\eta}(\mathbf{x}) \\ v_{\eta}(\mathbf{x}) \end{pmatrix}$

W Zwerger, Ed., *The BCS-BEC crossover and the Unitary Fermi Gas*, 2012. ✓ Bulgac, Yu., PRL 88, 2002.

#### Asymmetric Superfluid Local Density Approximation (ASLDA)

- 1. Extends the BdG functional to the Fermi superfluid gas at unitarity (UFG).
- 2. The functional reads as

$$\mathcal{E}(\mathbf{x}) = \frac{\hbar^2}{2m} \left( \alpha_{\uparrow}(n_{\uparrow}, n_{\downarrow})\tau_{\uparrow} + \alpha_{\downarrow}(n_{\uparrow}, n_{\downarrow})\tau_{\downarrow} + D(n_{\uparrow}, n_{\downarrow}) \right) + \frac{|\Delta|^2}{g(n_{\uparrow}, n_{\downarrow})}$$

with parameters  $\alpha_{\uparrow}$  and D have to be determined ab initio (e.g. Quantum Monte Carlo simulations).

Here  $\mathcal{H}_{BdG} = -\hbar^2 \nabla^2 / 2m$ 

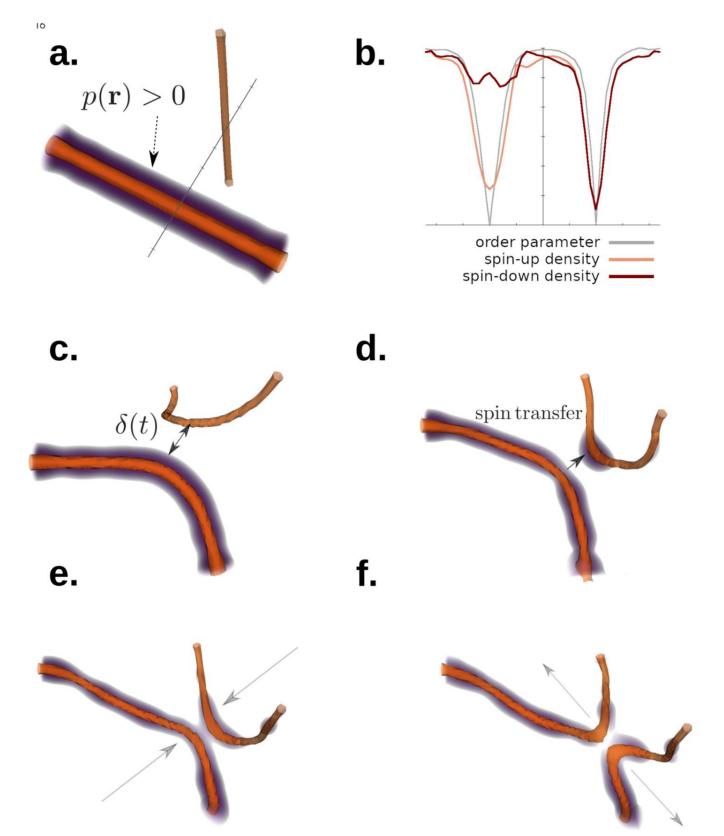
3. The renormalization requires the energy cut-off and a renormalized  $g_{eff}$ 

3. The minimization procedure consists of self-consistent iterations of the modified BdG equations similarly to the BCS case.

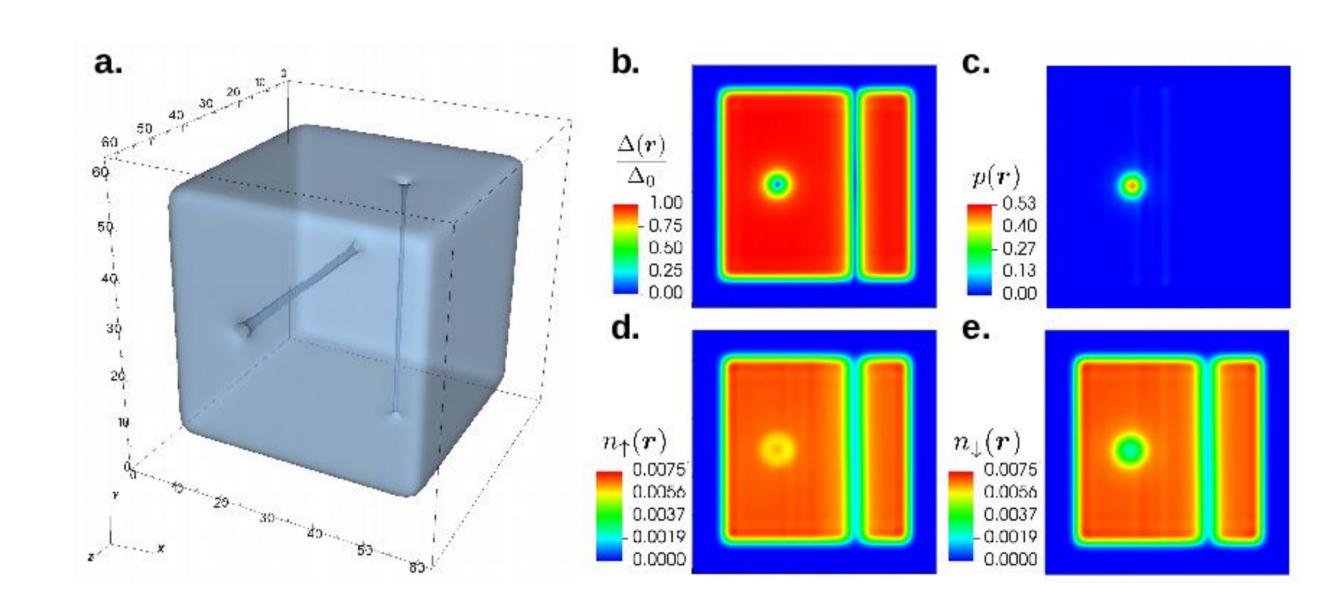
4. By replacing  $E_{\eta} \rightarrow i\hbar \frac{\partial}{\partial t}$  we obtain time-dependent version of ASLDA equations.

# 2. Vortex Collisions with a spin polarization

- 1. Initially perpendicular vortex lines, where only one carries a spin polarization.
- 2. At short distances vortex lines interact due to phase configuration.
- 3. Before the reconnection event there is a spin polarization transfer between the approaching vortex lines.
- 4. The vortex lines tend to assume an antiparallel configuration before they reconnect.
- 5. After the reconnection the two vortex lines have a similar spin polarization.



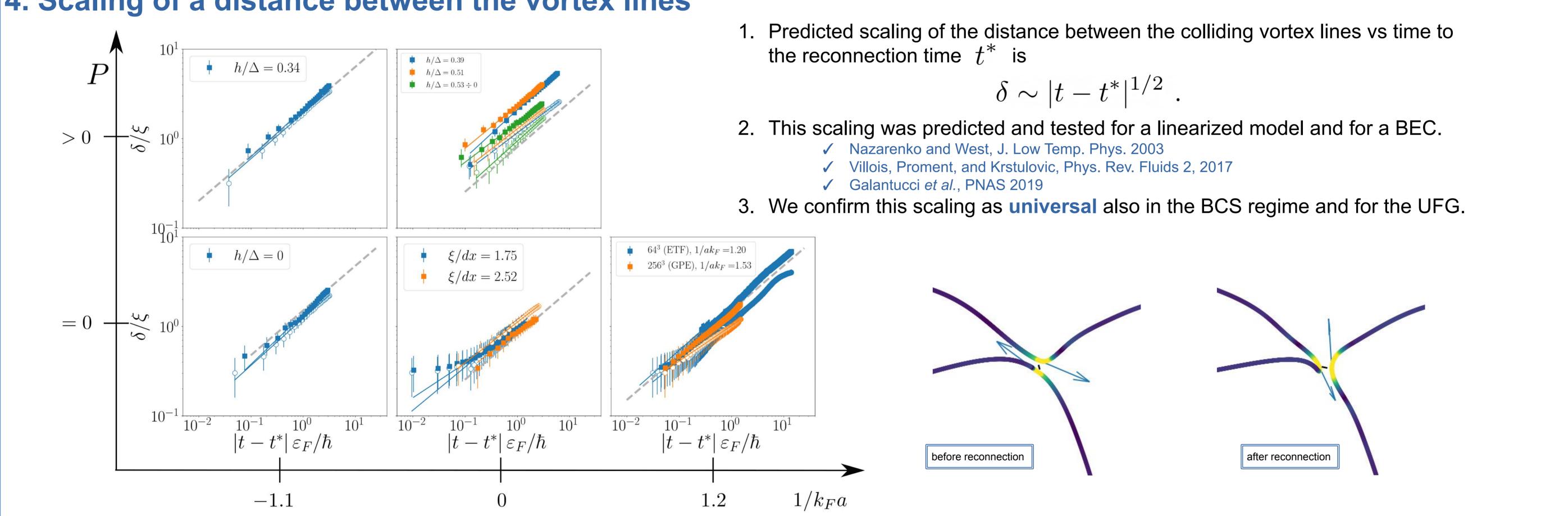
# 3. Initial and boundary conditions



Initially vortex lines are perpendicular.

- 2. At short distances vortex lines interact mostly due to phase configuration.
- The potential wall at the boundary is chosen such that it is smooth 3. (vanishing all derivatives).

#### 4. Scaling of a distance between the vortex lines



## **5. High Performance Computing**

Simulations methods

Future plans

- ✓ Simulations are conducted on a 3D spatial lattice of size  $N_x \times N_y \times N_z$  with no symmetry restrictions.
- ✓ We use massively parallelized numerical codes with GPU acceleration
- ✓ With present supercomputers capabilities we can study 3D dynamics of systems containing up to  $3 \times 10^5$  of ultracold fermions, with arbitrary spin-polarization.

#### Selected hpc systems in use by our group



Piz Daint @ Swiss National Summit @ Oak Ridge National Supercomputing Centre (Switzerland) Laboratory (USA)

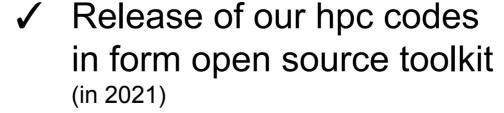




and Computing Center, (Tokyo)

Tsubame3.0 @ Global Scientific Information Prometheus @ Academic Computer

Centre CYFRONET (Poland)



Extension of the toolkit towards Bose-Fermi mixtures

Acknowledgments

PRACE

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