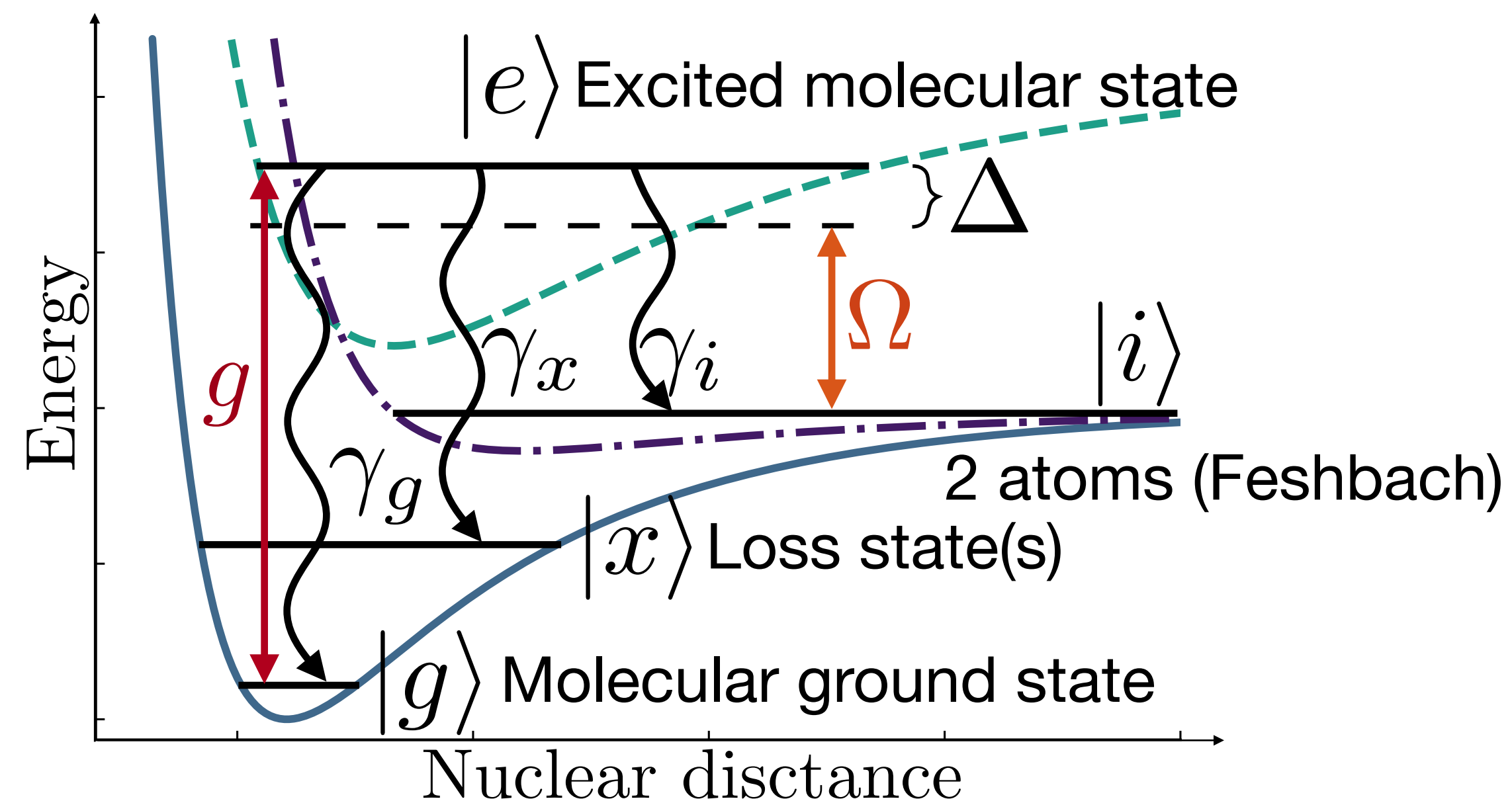
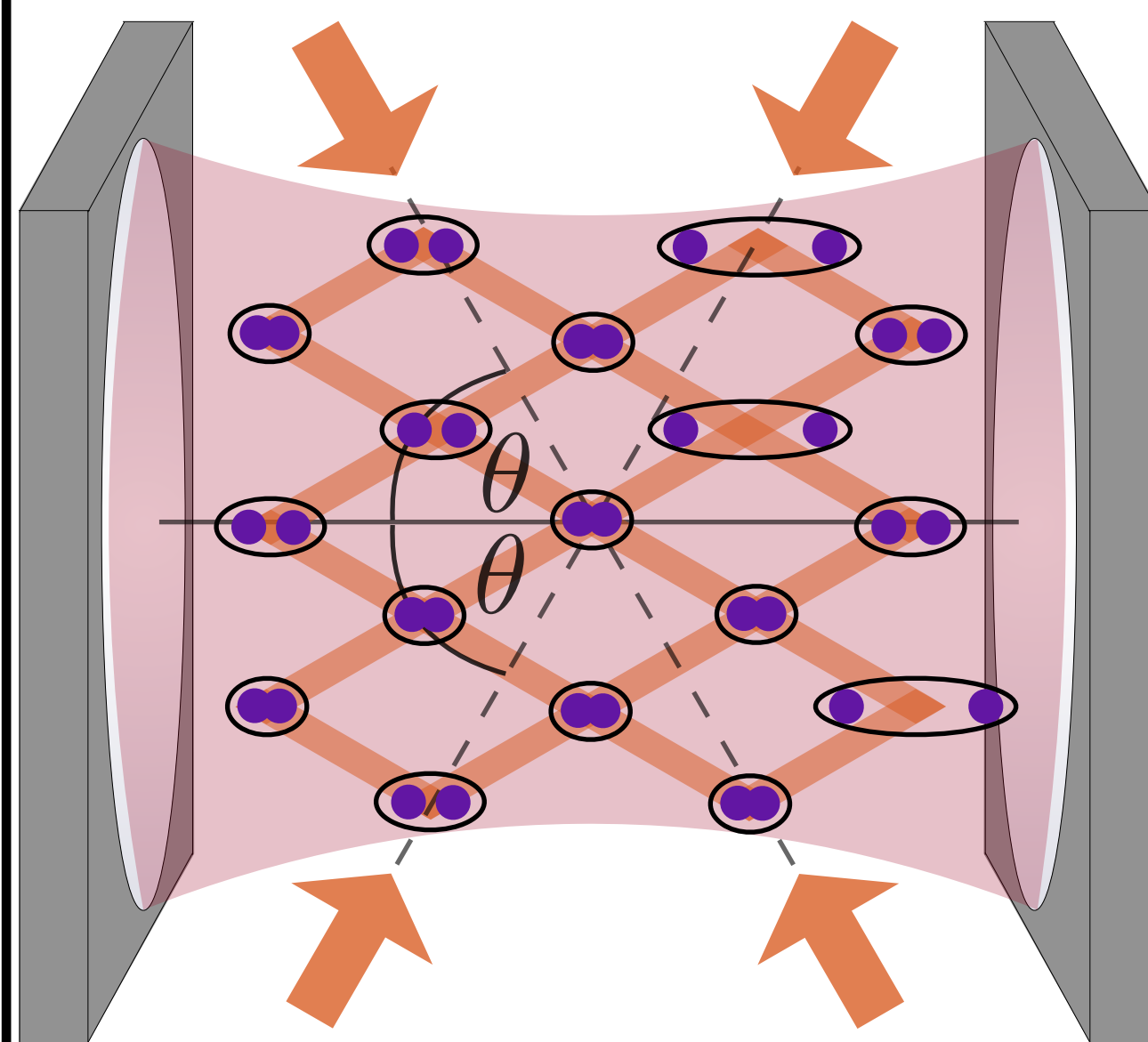


# Abstract / Motivation:

Ultracold atoms and molecules that are coupled to a cavity can exhibit collective effects, such as polariton formation or superradiance. Here, we propose a mechanism to harness these collective and dissipative effects in order to realize high-yield ground state molecular formation from ultracold atoms. Using a combination of analytical and numerical techniques, we demonstrate that the ground state yield can be improved by increasing the number of atoms, and overcome efficiencies of state-of-the-art association schemes.

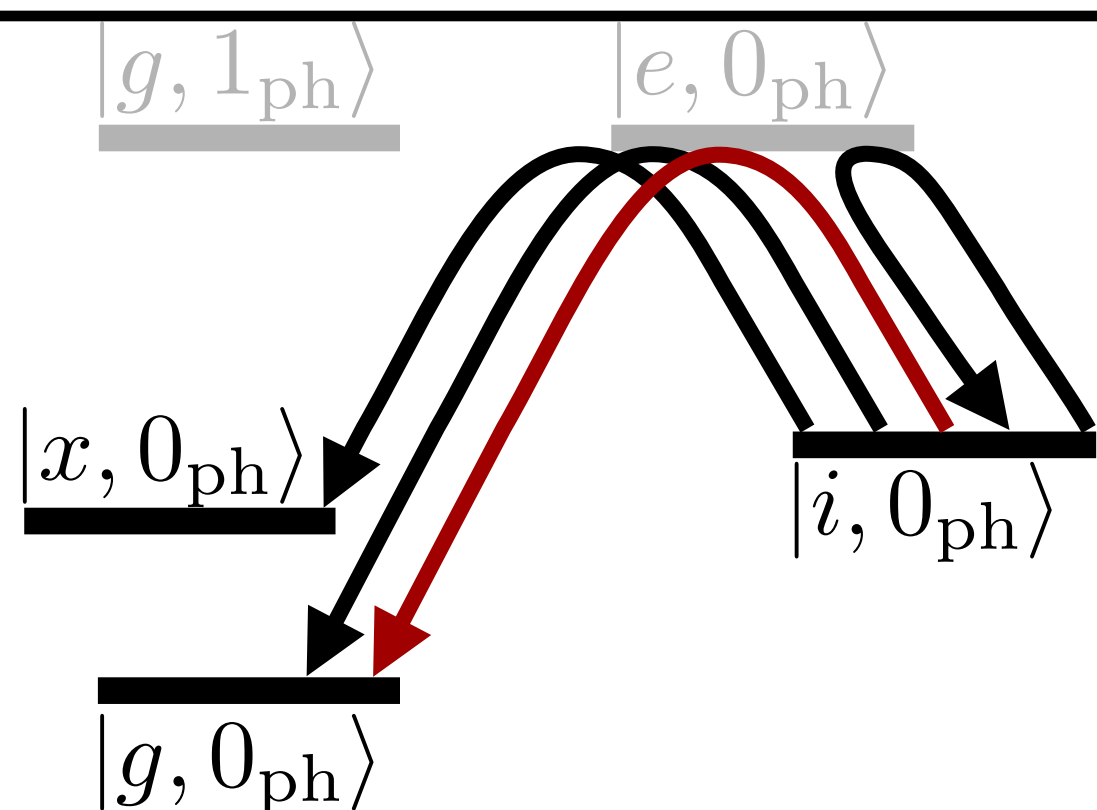
# Model

DW, Schütz, Whitlock, Schachenmayer, Pupillo: PRL 125, 193201 (2020)

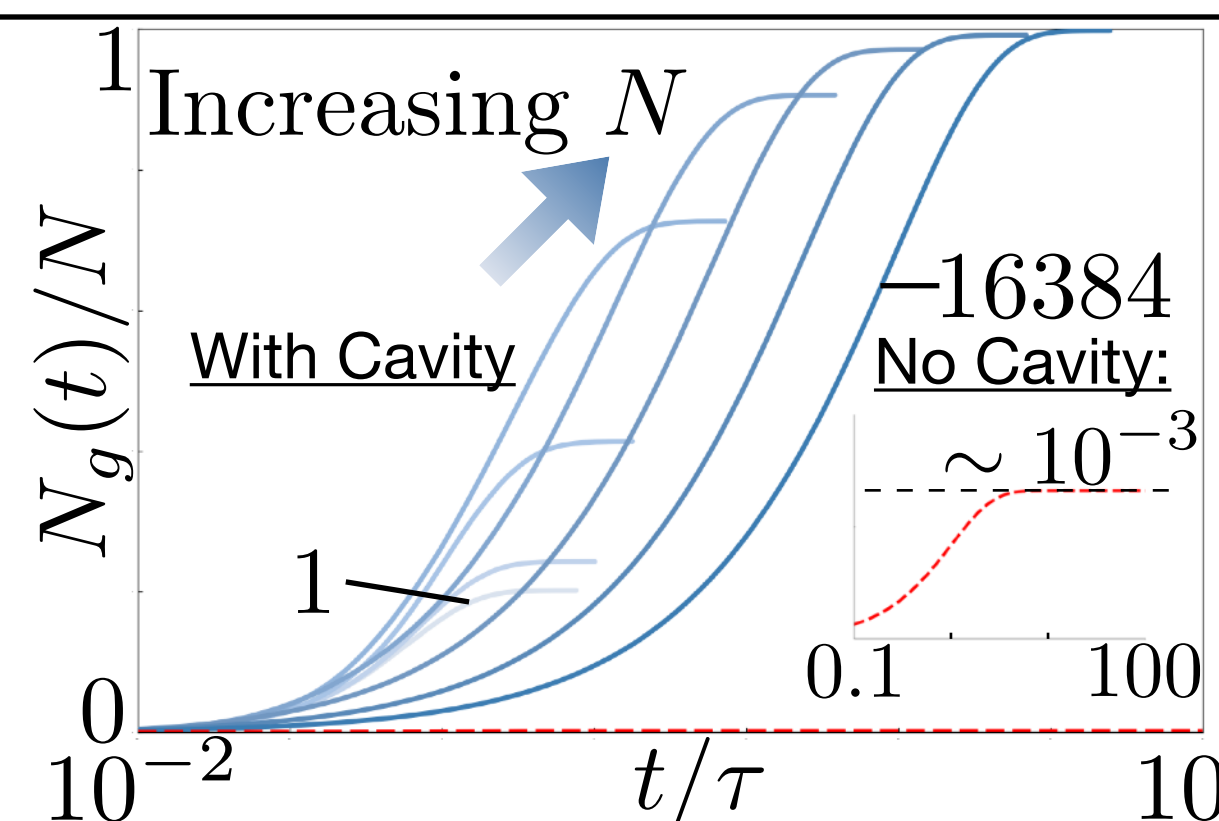


See also: Kampschulte, Hecker Denschlag: NJP 20, 123015 (2018)

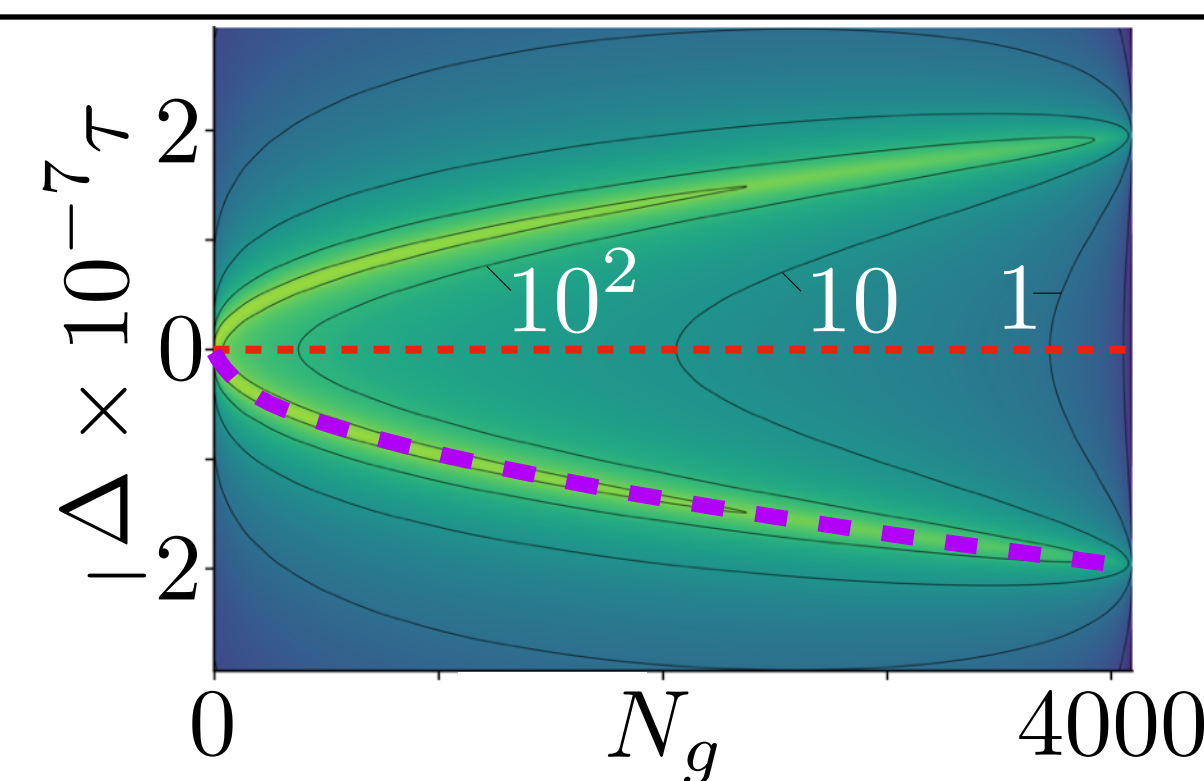
Perez-Rios, Kim, Hung: NJP 19, 123035 (2017)



# Adiabatic Elimination



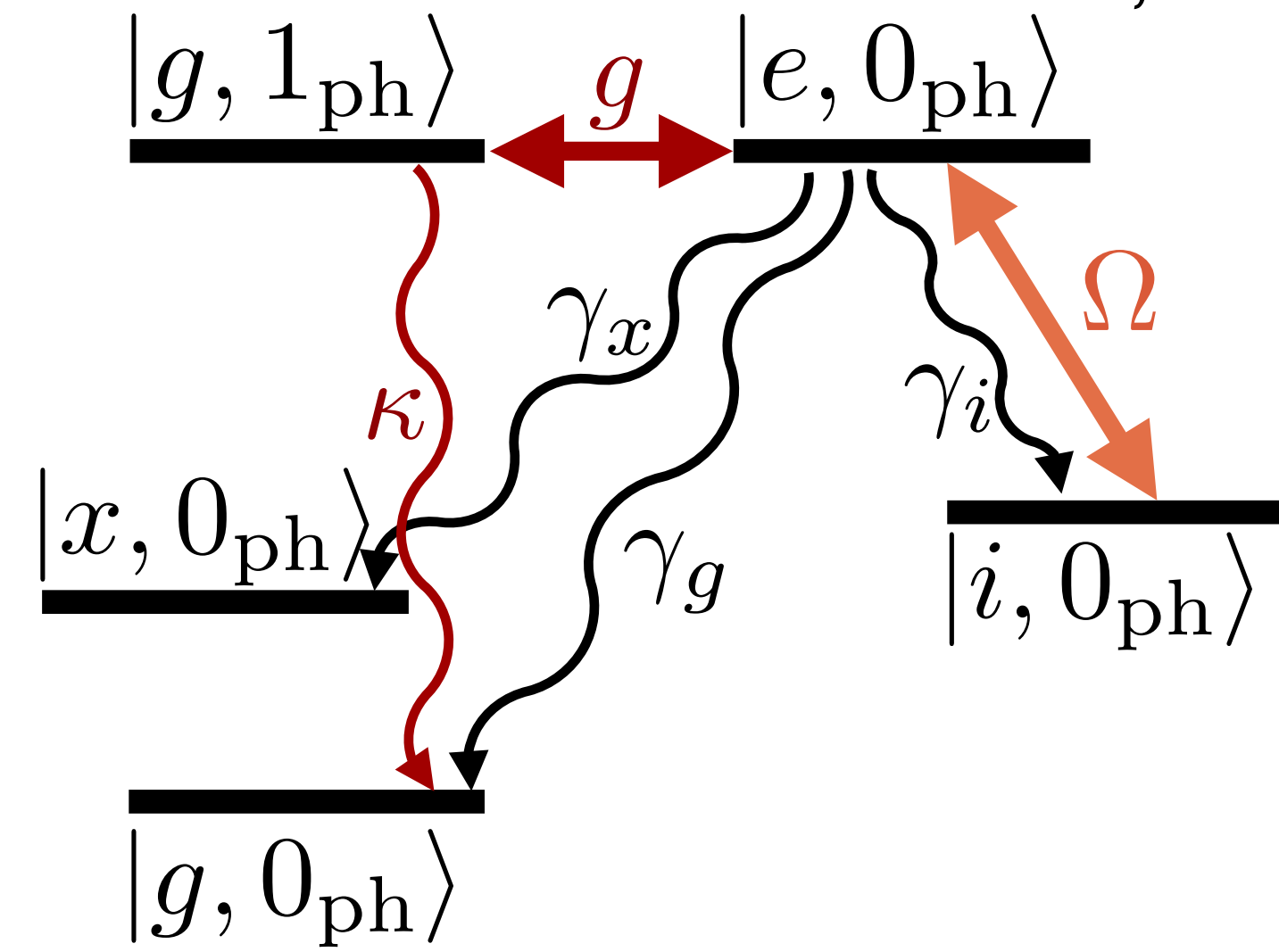
# Collective Enhancement



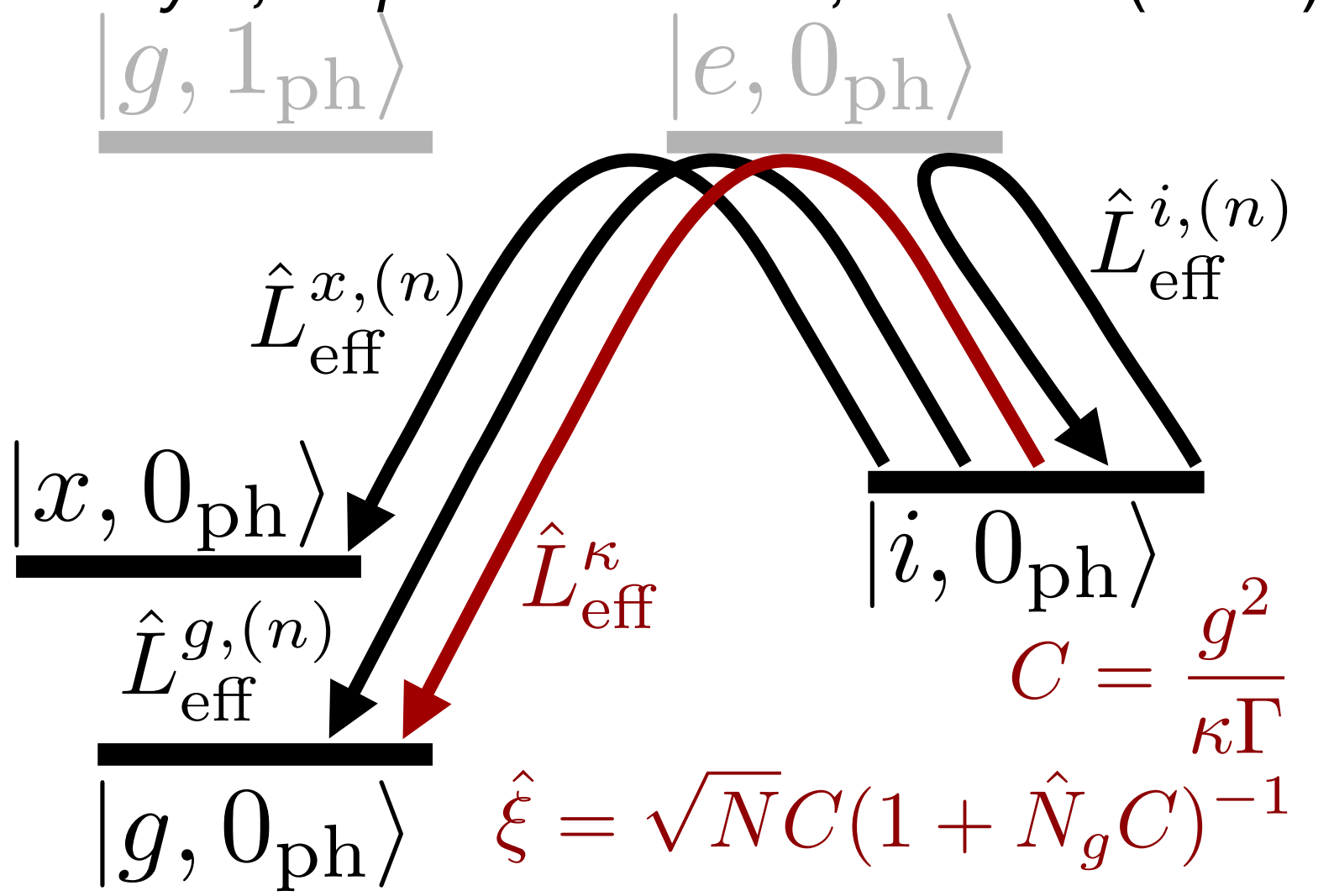
# Polaritons & Rate Equations

# Method

The full dynamics can only be simulated for small systems. In the weak driving limit, the excited state and cavity mode have a small population and can be adiabatically eliminated. The effective dynamics is purely dissipative and can be efficiently simulated using the permutation symmetry.



Eliminate for small  $\Omega$   
**weak driving limit**



**Master Equation**

$$\partial_t \hat{\rho} = -i[\hat{H}, \hat{\rho}] + \sum_{\alpha,n} \mathcal{L}_{\alpha,(n)}(\hat{\rho}) + \mathcal{L}_{\kappa}(\hat{\rho})$$

$$\mathcal{L}_k(\hat{\rho}) = -\hat{L}^{k\dagger} \hat{L}^k \hat{\rho} - \hat{\rho} \hat{L}^{k\dagger} \hat{L}^k + 2\hat{L}^k \hat{\rho} \hat{L}^{k\dagger}$$

**Hamiltonian and Lindblad**

$$\hat{H}_{LA} = \Omega\sqrt{N} (\hat{S}_{ie} + \hat{S}_{ei}) + \Delta\hat{N}_e$$

$$\hat{H}_C = g\sqrt{N} (\hat{a}^\dagger \hat{S}_{ge} + \hat{S}_{eg} \hat{a}) + \Delta\hat{a}^\dagger \hat{a}$$

$$\hat{L}_{\kappa} = \sqrt{\kappa} \hat{a} \quad \hat{L}_{\alpha}^{(n)} = \sqrt{\gamma_{\alpha}} \hat{\sigma}_{\alpha e}^{(n)}$$

**Effective Hamiltonian and Lindblad**

$$\hat{H}_{eff} = 0$$

$$\hat{L}_{eff}^{\kappa} = \frac{\Omega\sqrt{\kappa}}{g} \hat{\xi} \hat{S}_{gi}$$

$$\hat{L}_{eff}^{\alpha,(n)} = \frac{\Omega\sqrt{\gamma_{\alpha}}}{\Gamma} \left( \hat{\sigma}_{\alpha i}^{(n)} - \hat{\sigma}_{\alpha g}^{(n)} \hat{\xi} \hat{S}_{gi} \right)$$

**Effective dynamics is purely dissipative and collective!**

Formalism: Reiter, Sørensen, PRA 85 032111 (2012)

Large scale quantum trajectories simulation possible using **permutation symmetry**

Symmetry: Chase, Geremia, PRA 78.5 052101 (2008) / Zhang, Zhang, Mølmer, NJP 20 112001 (2018)

$$\hat{\sigma}_{\alpha,\beta}^{(n)} = |\alpha\rangle \langle \beta|_n$$

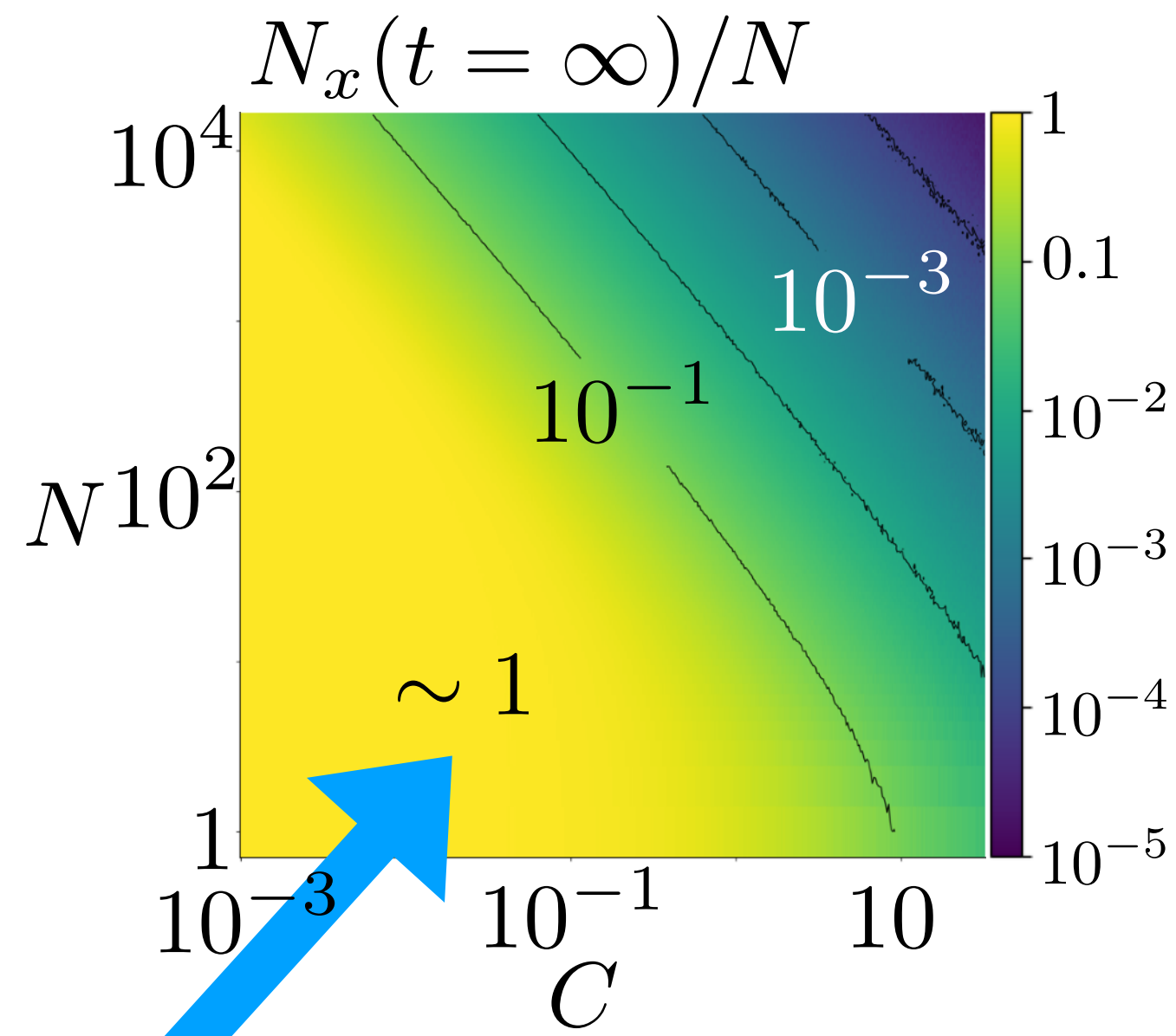
$$\hat{N}_{\alpha} = \sum \hat{\sigma}_{\alpha\alpha}^{(n)}$$

$$\hat{S}_{\alpha\beta} = \sum_n \sigma_{\alpha\beta}^{(n)} / \sqrt{N}$$



# Collective Enhancement

The cavity coupling leads to a collective enhancement of the final ground state population, and a collective slowdown. The decay is mediated by the cavity, which transfers population to the ground state, and induces a slowdown due to polariton or Zeno blocking. This can also be seen from the effective rate equations (next slide). The relevant scaling parameter is the **collective cooperativity**  $NC$ .



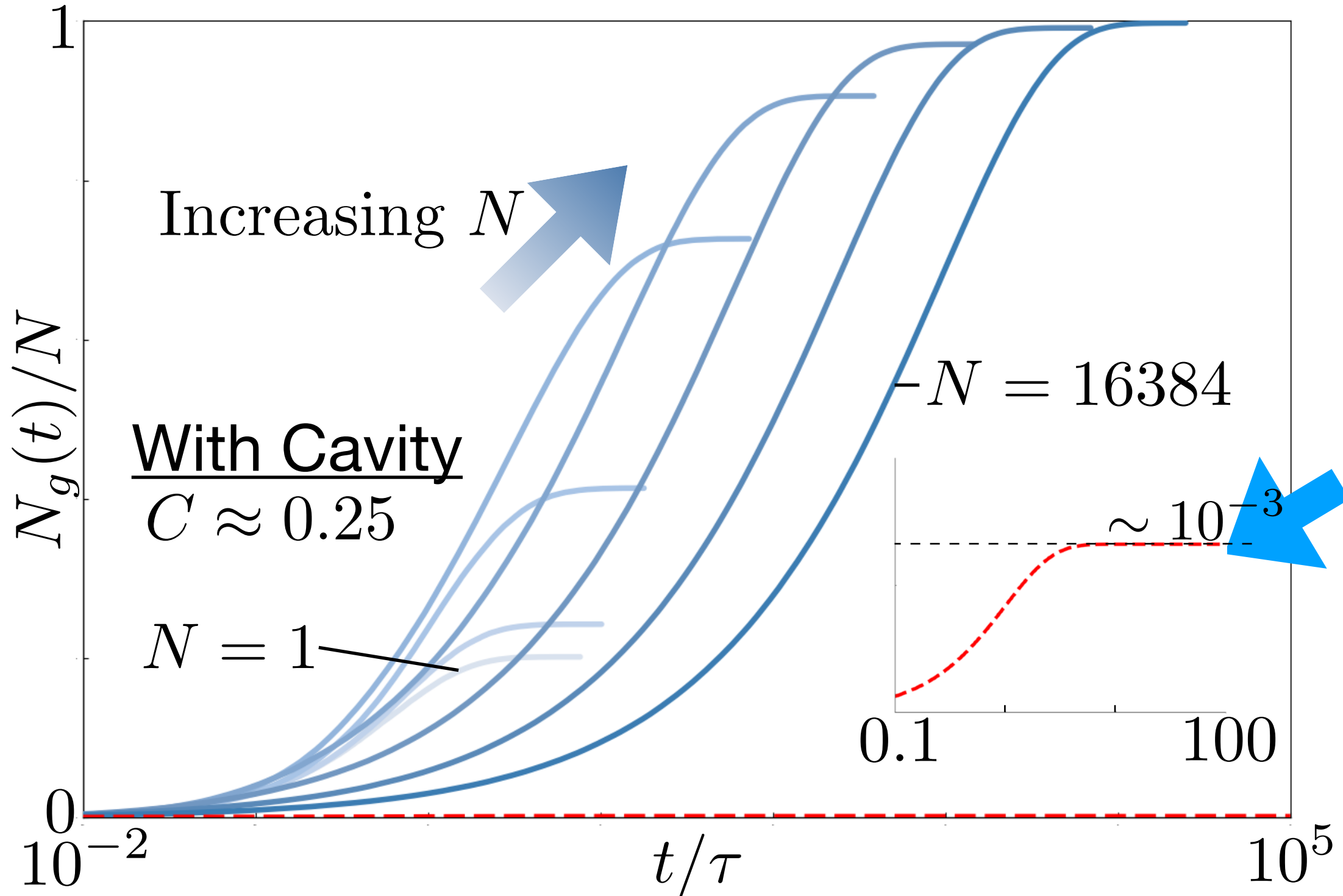
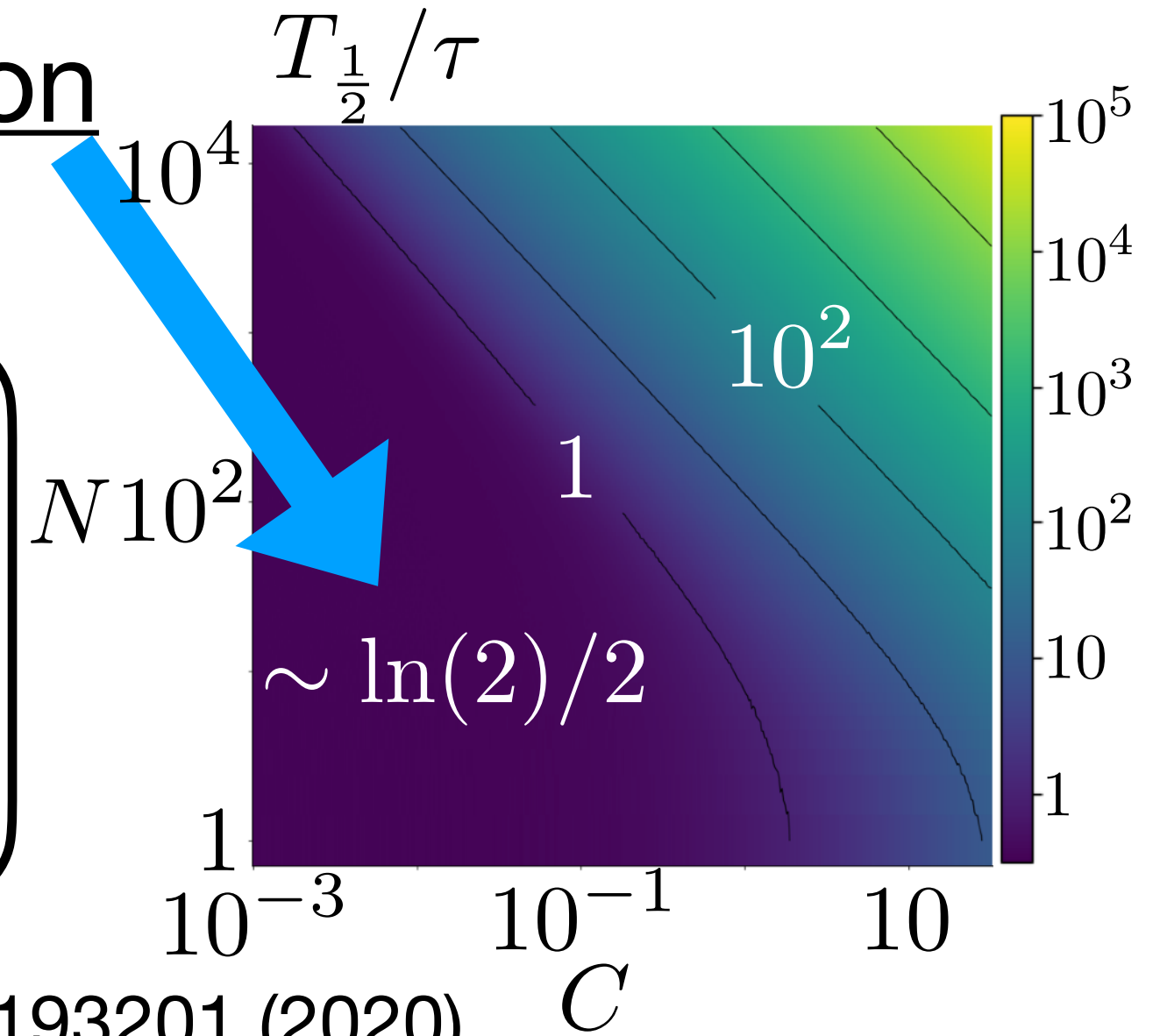
**Increased efficiency:**

- Superradiance
- Cavity mediated decay

**No Cavity: Photoassociation**

**Slowdown:**

- Zeno blocking
- Polariton detuning



Characteristic Timescale  $\tau = \Gamma/\Omega^2$

Cavity Cooperativity  $C = \frac{g^2}{\kappa\Gamma}$

# Rate Equations

Rate equations for state populations:

$$\dot{N}_i \approx -\frac{2}{\tau} \frac{N_i}{(N_g + 1)C} \approx -\dot{N}_g$$

$$\dot{N}_x \approx \frac{2f_x}{\tau} \frac{N_i}{(N_g + 1)^2 C^2} \ll \dot{N}_g$$

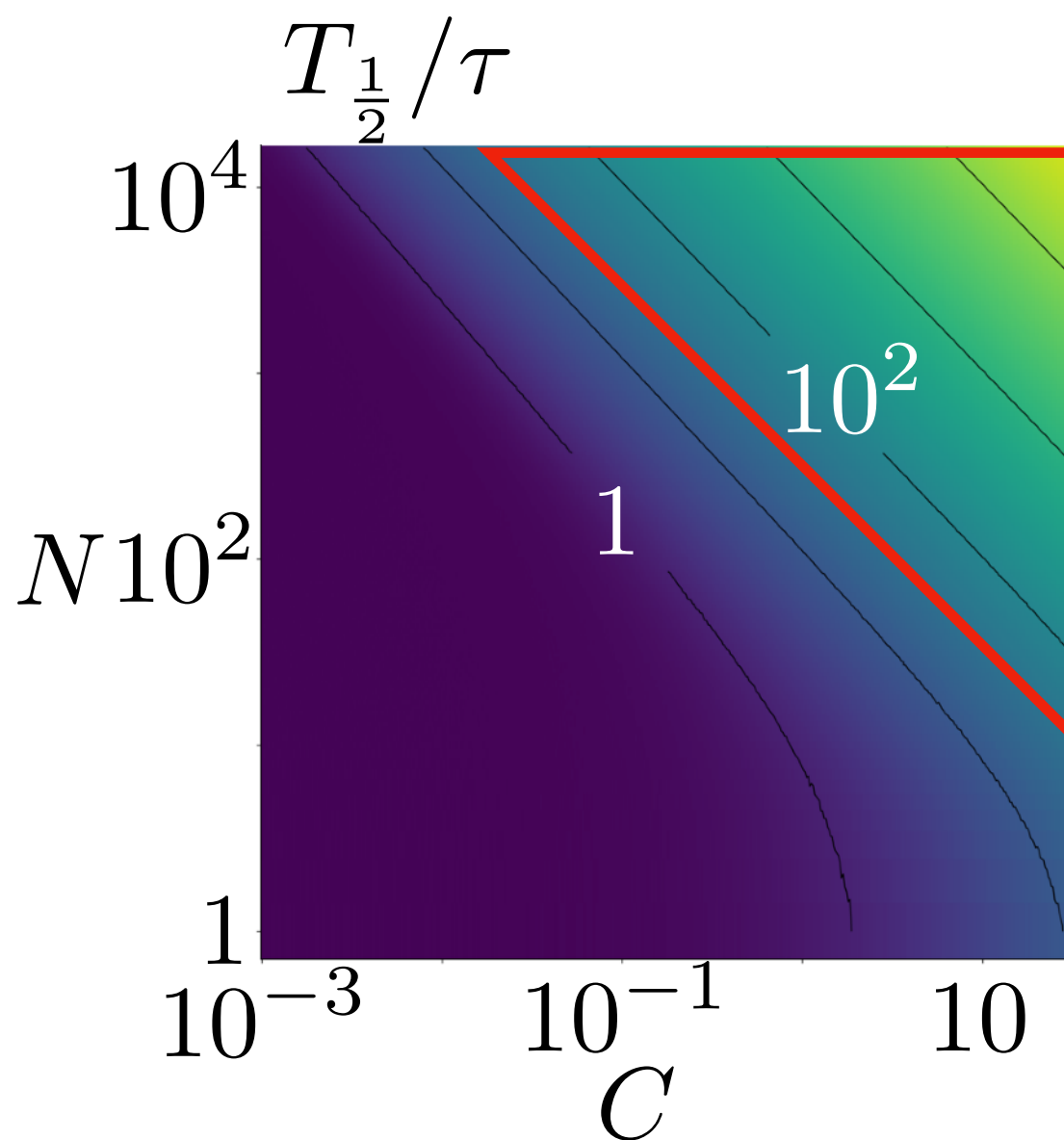
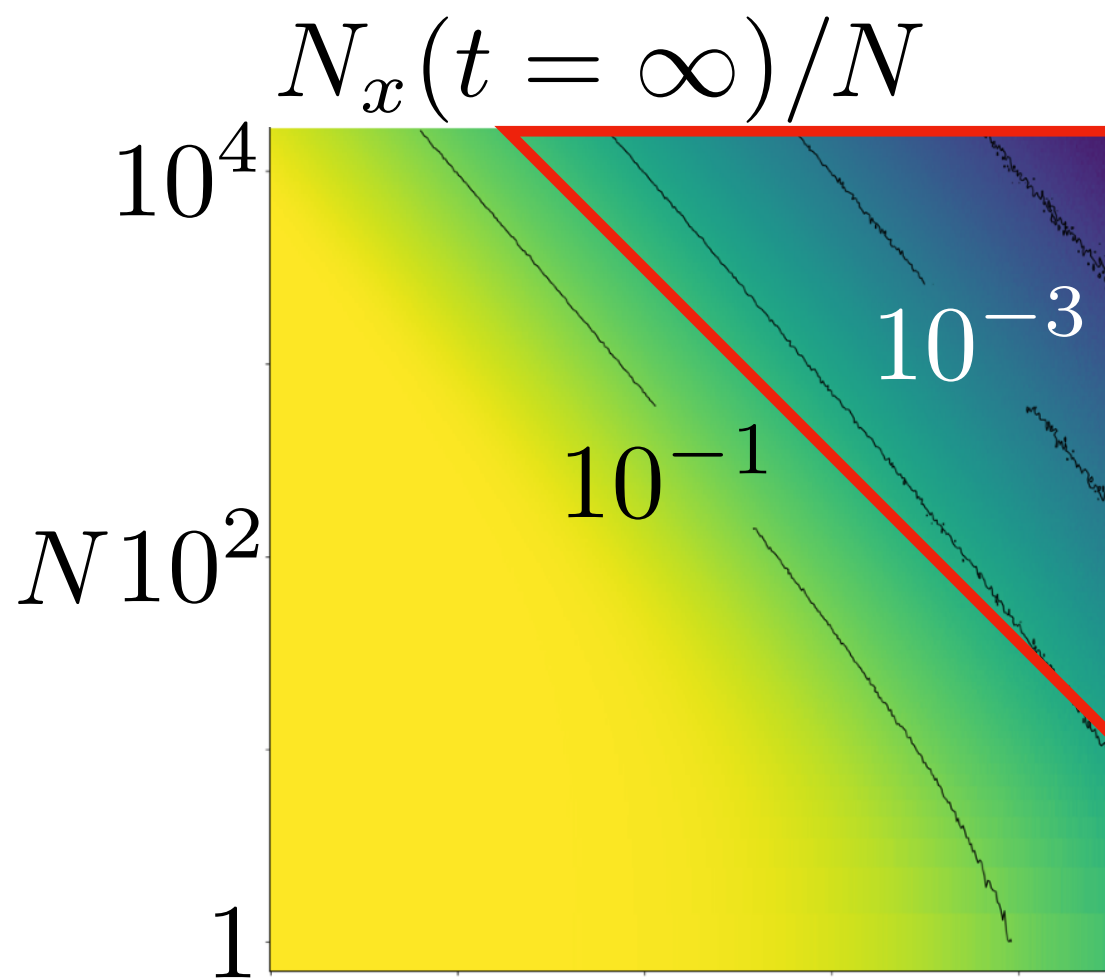
↓ solve

**Final loss state population**

$$\frac{N_x(t = \infty)}{N} \approx \frac{f_x \ln(N)}{NC}$$

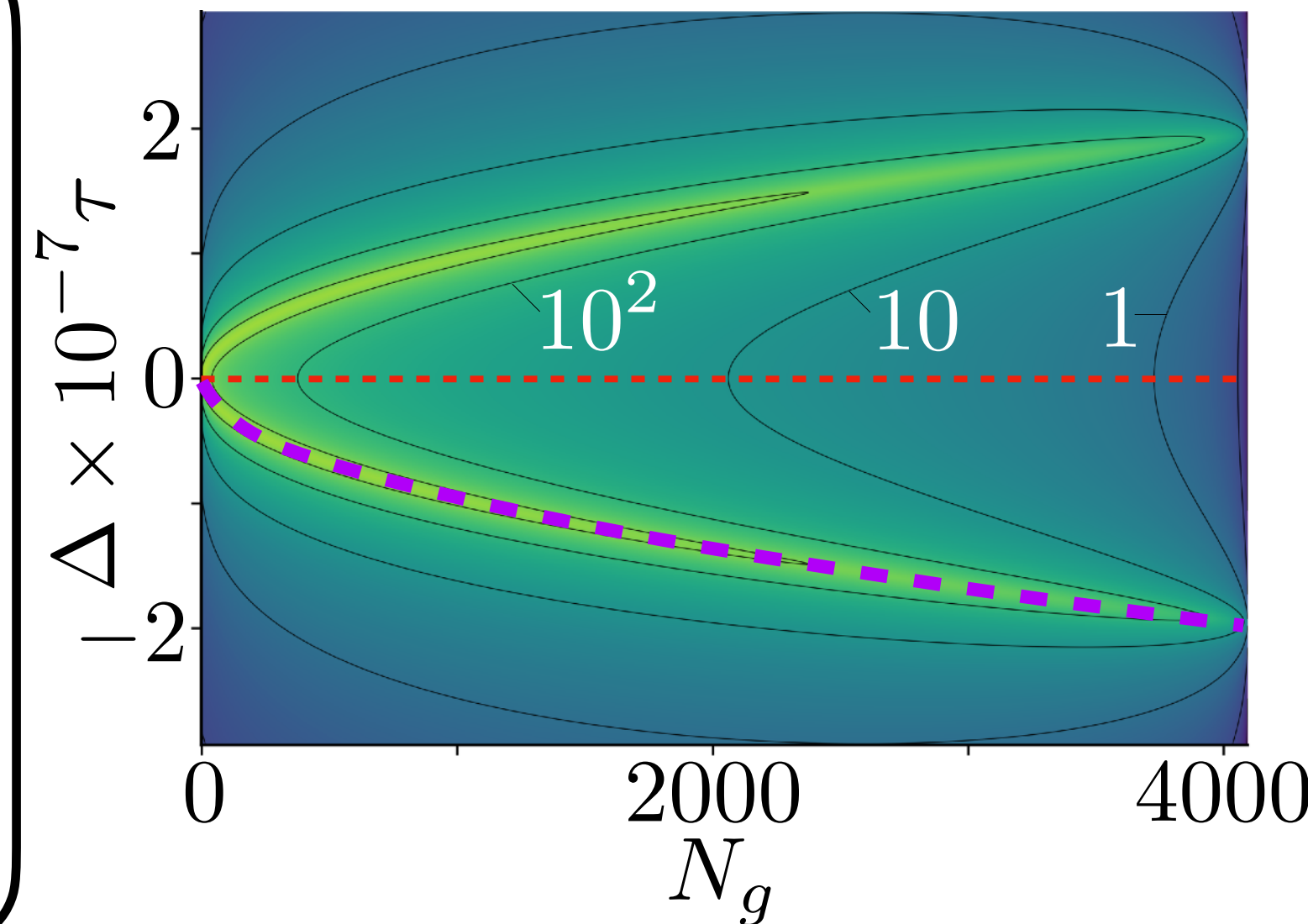
**Initial state half time**

$$T_{\frac{1}{2}} \approx NC\tau \left( \ln(2) - \frac{1}{2} \right)$$



# Slowdown $\Delta(t)$

The slowdown caused by the detuning from the virtual polariton resonance can be reduced by chirping the laser. This leads to an efficiency and a life time that scale with  $\kappa$ .



**Rate equations and solutions**

$$\dot{N}_i = -\frac{2\Omega^2(\kappa + \gamma_g + \gamma_x)}{(\kappa + \Gamma)^2} N_i \quad \dot{N}_x = -\frac{2\Omega^2\gamma_x}{(\kappa + \Gamma)^2} N_i$$

$$T_{\frac{1}{e}} \approx \frac{(\kappa + \Gamma)^2}{2\Omega^2(\kappa + \gamma_g + \gamma_x)} \quad \frac{N_x(t = \infty)}{N} \approx \frac{\gamma_x}{\kappa + \gamma_g + \gamma_x}$$

valid for  $g \ll \kappa \ll \sqrt{Ng}$

Slowdown can be reduced by staying on **polariton resonance**